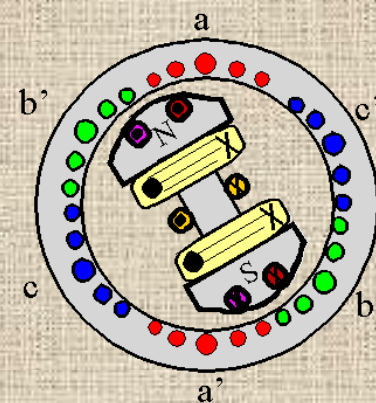
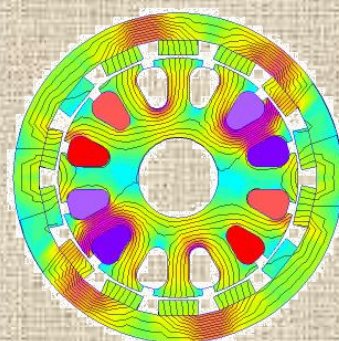
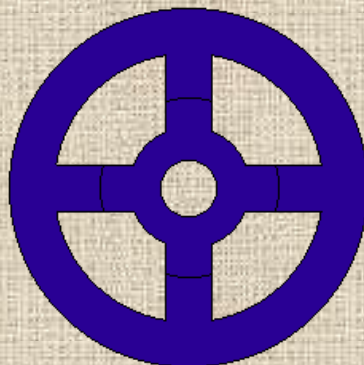
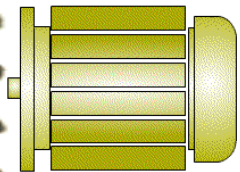


# EE552 ELECTRICAL MACHINES III

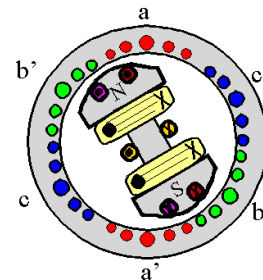


## LECTURE 10



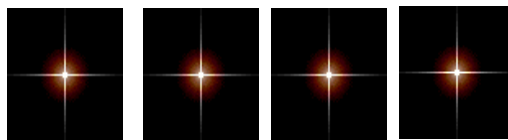
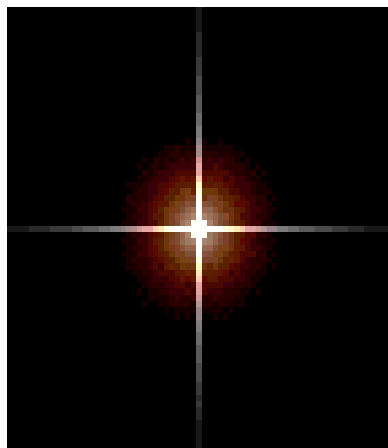


# LECTURE NOTES



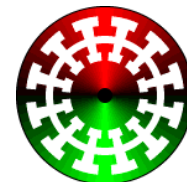
## ELECTRICAL MACHINES III

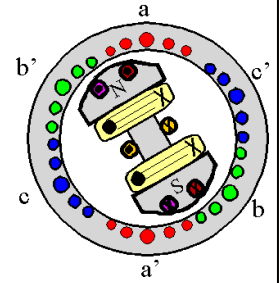
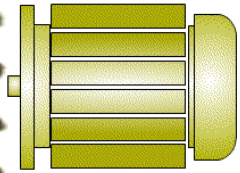
EE552



SPRING 2018

Dr : MUSTAFA AL-REFAI

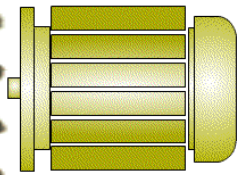




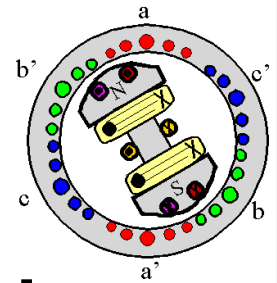
# LECTURE 10

## SYNCHRONOUS GENERATOR





# Rotation speed of synchronous generator



**By the definition, synchronous generators produce electricity whose frequency is synchronized with the mechanical rotational speed.**

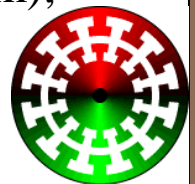
$$f_e = \frac{n_m P}{120}$$

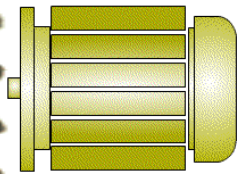
Where  $f_e$  is the electrical frequency, Hz;

$n_m$  is mechanical speed of magnetic field (rotor speed for synchronous machine), rpm;

$P$  is the number of poles.

Steam turbines are most efficient when rotating at high speed; therefore, to generate 60 Hz, they are usually rotating at 3600 rpm and turn 2-pole generators. Water turbines are most efficient when rotating at low speeds (200-300 rpm); therefore, they usually turn generators with many poles.





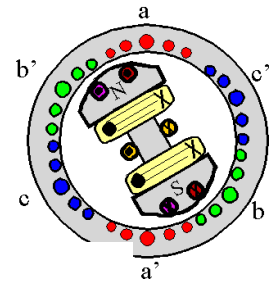
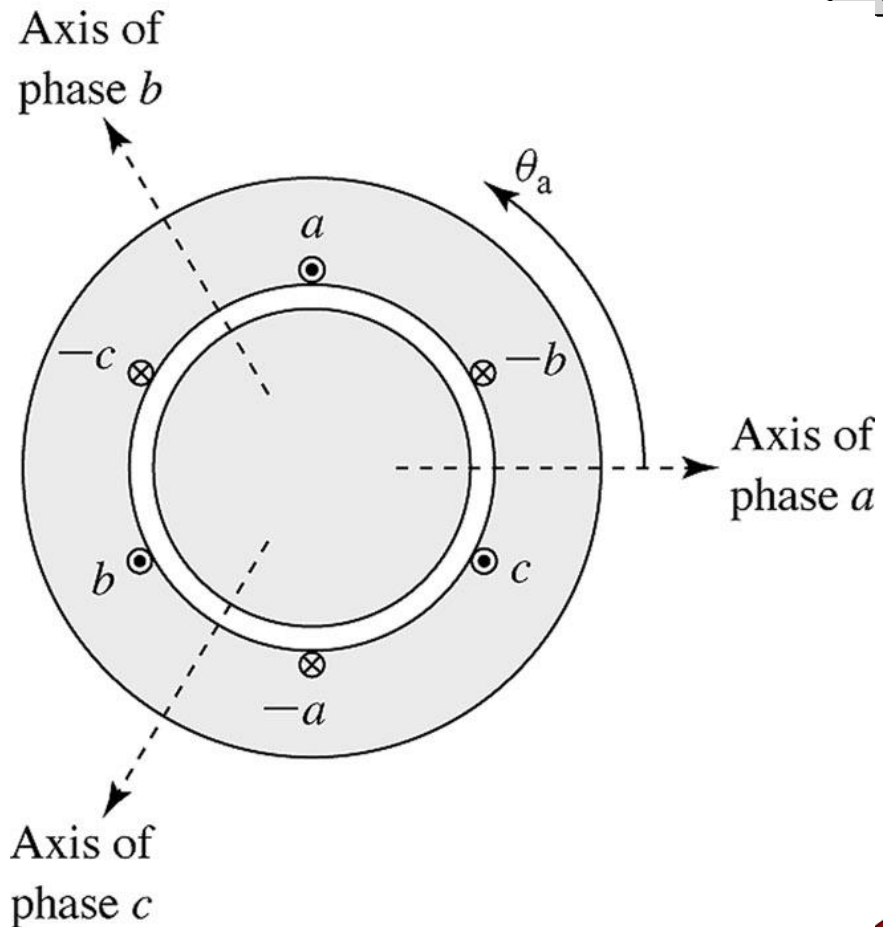
# MMF Wave of a Polyphase Winding

**Simplified two-pole three-phase stator winding.**

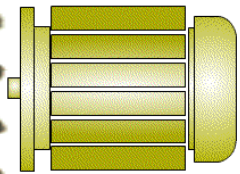
$$i_a = I_m \cos \omega_e t$$

$$i_b = I_m \cos(\omega_e t - 120^\circ)$$

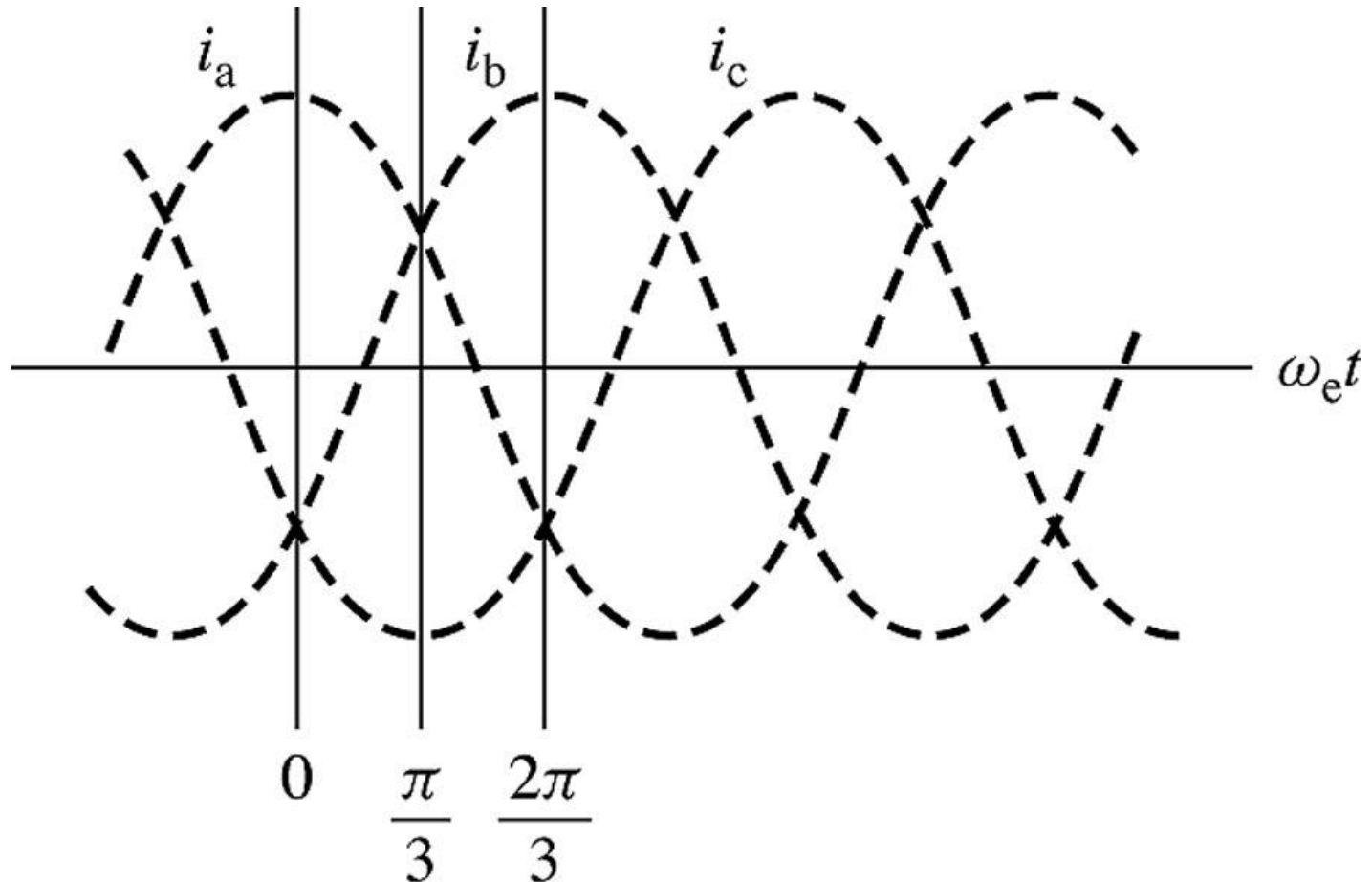
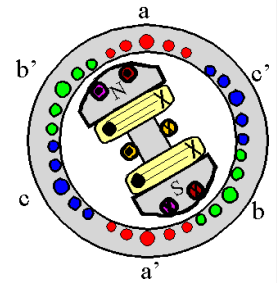
$$i_c = I_m \cos(\omega_e t + 120^\circ)$$

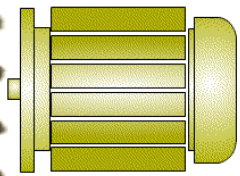






# Instantaneous phase currents under balanced three-phase conditions.

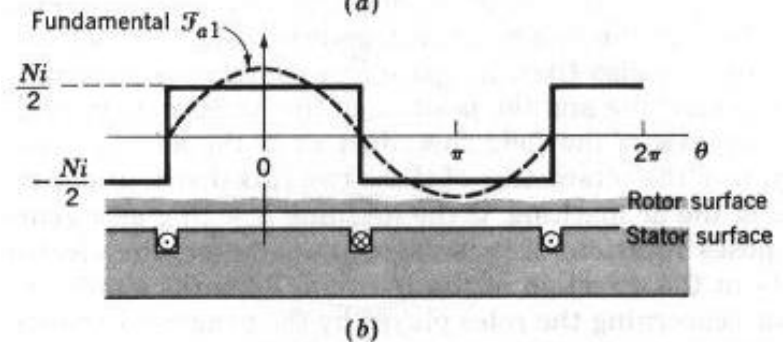
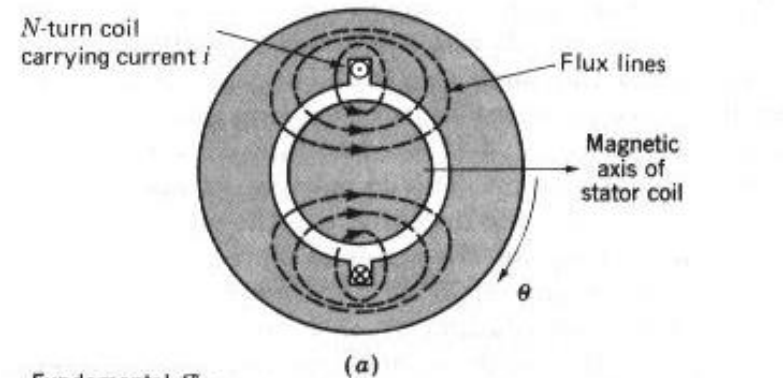
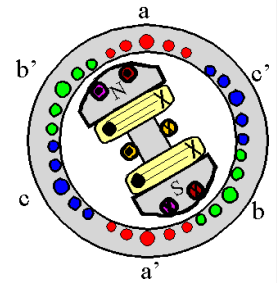




# Magnetomotive Force (mmf) of AC Windings

## M.m.f. of a coil

- the variation of magnetic potential difference along the air-gap periphery is of rectangular waveform and of magnitude  $\frac{1}{2}Ni$
- The amplitude of mmf wave varies with time, but not with space
- The air-gap mmf wave is time-variant but space invariant
- The air-gap mmf wave at any instant is rectangular



The mmf of a concentrated full-pitch coil.

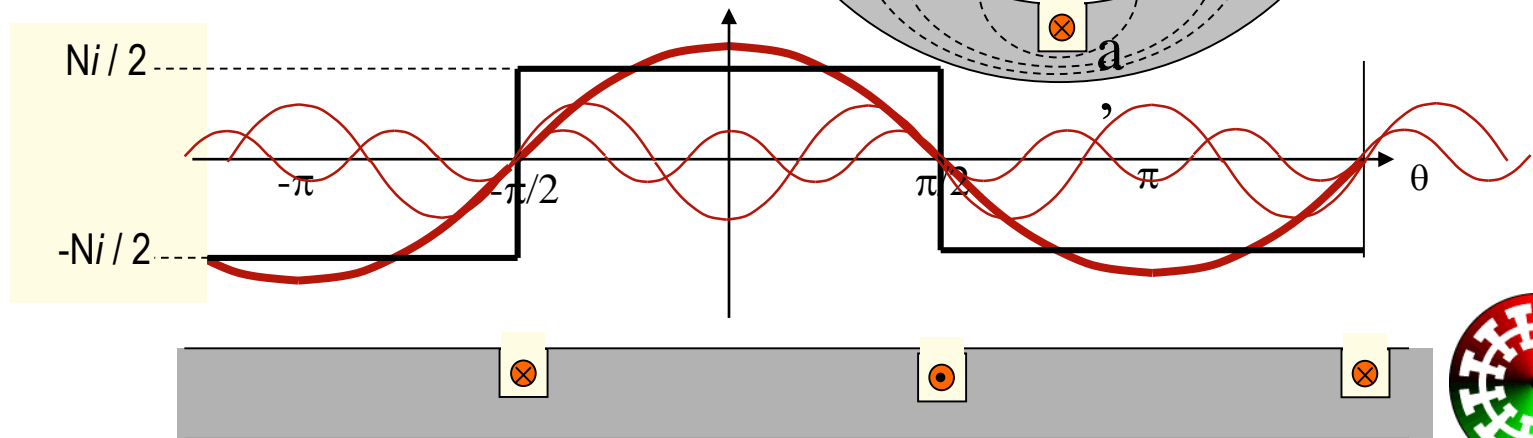
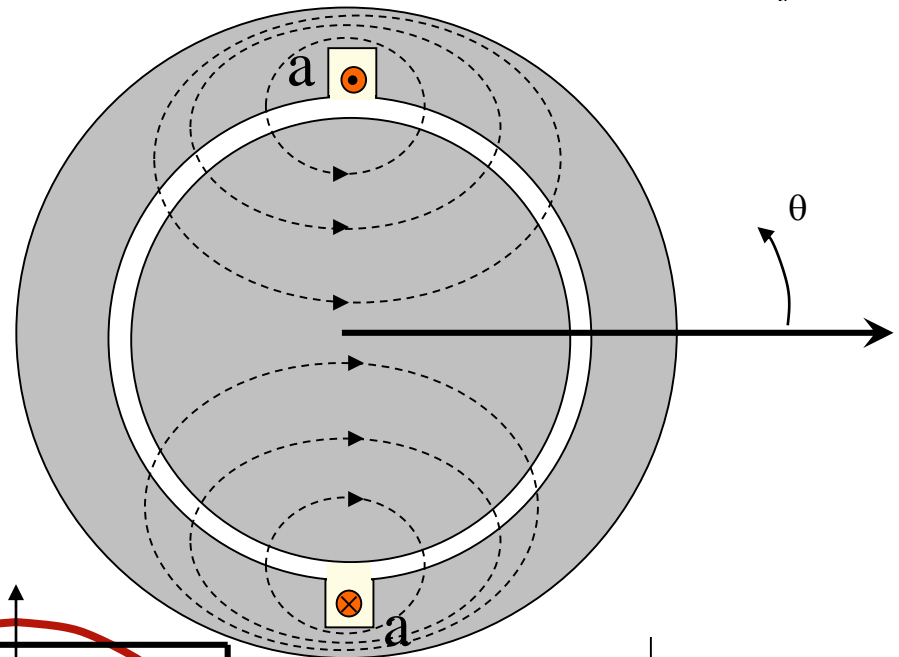
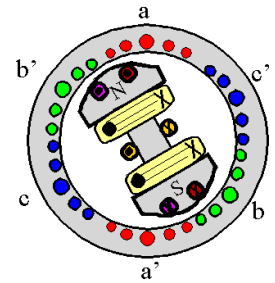


# Rotating Magnetic Fields

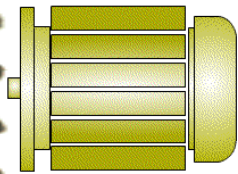
□ Single N turn coil carrying current  $i$

Spans  $180^\circ$  elec

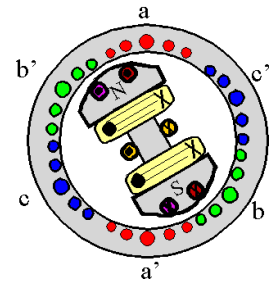
□ Permeability of iron  $\gg \mu_0$   
 $\rightarrow$  all MMF drop appear in airgap







# Rotating Magnetic Fields cont...



□ The fundamental component of rectangular wave is found to be

$$F_{a1} = \frac{4}{\pi} \cdot \frac{Ni}{2} \cos \alpha = F_{1p} \cos \alpha$$

Where

$\alpha$  = electrical space angle measured from the magnetic axis of the stator coil

Here  $F_{1p}$ , the peak value of the sine mmf wave for a 2-pole machine is given by

$$F_{1p} = \frac{4}{\pi} \cdot \frac{Ni}{2} \text{ AT per pole}$$

When  $i=0 \rightarrow F_{1p}=0$

$$i = I_{max} = \sqrt{2}I$$

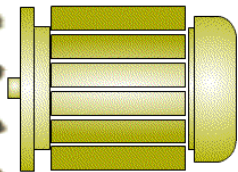
For 2-pole machine

$$F_{1pm} = \frac{4}{\pi} \cdot \frac{N\sqrt{2}I}{2} \text{ AT per pole}$$

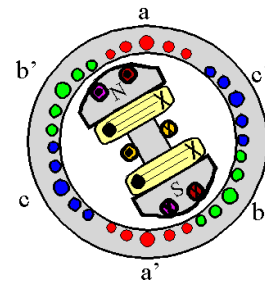
For p-pole machine

$$F_{1pm} = \frac{4}{\pi} \cdot \frac{N\sqrt{2}I}{p} \text{ AT per pole}$$



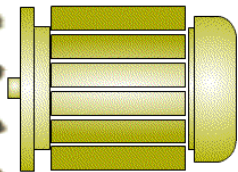


# M.M.F of Distributed Winding

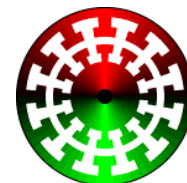
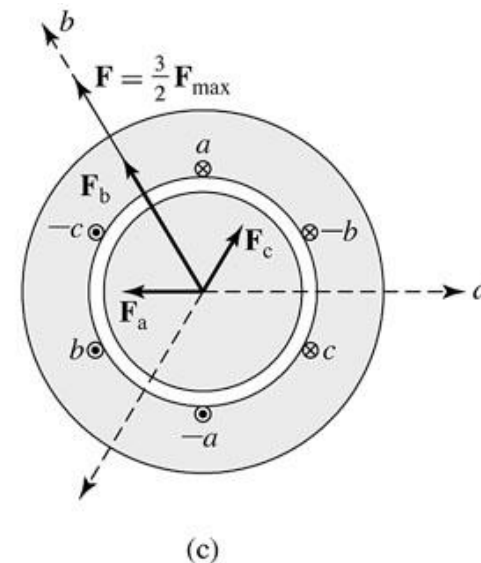
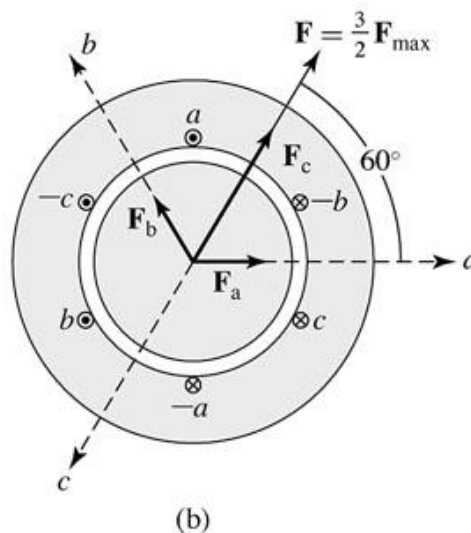
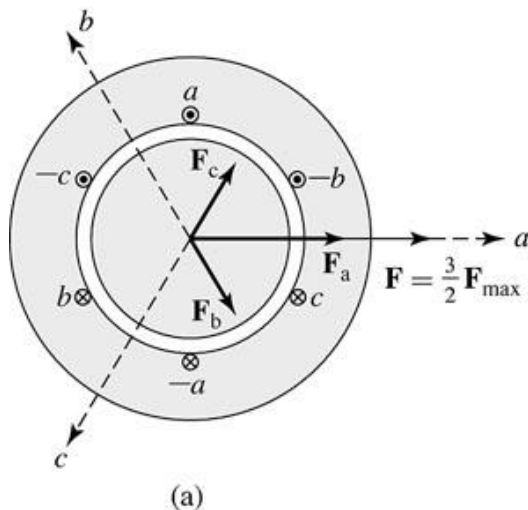
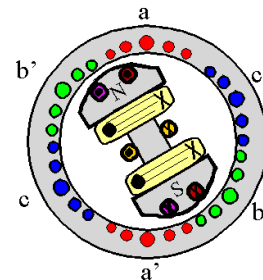


- ❑ The mmf distribution along the air gap periphery depends on the nature of slots, winding and the exciting current
- ❑ The effect of winding distribution has changed the shape of the mmf wave, from rectangular to stepped

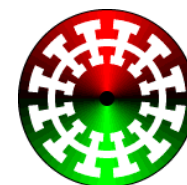
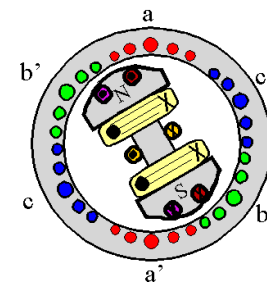
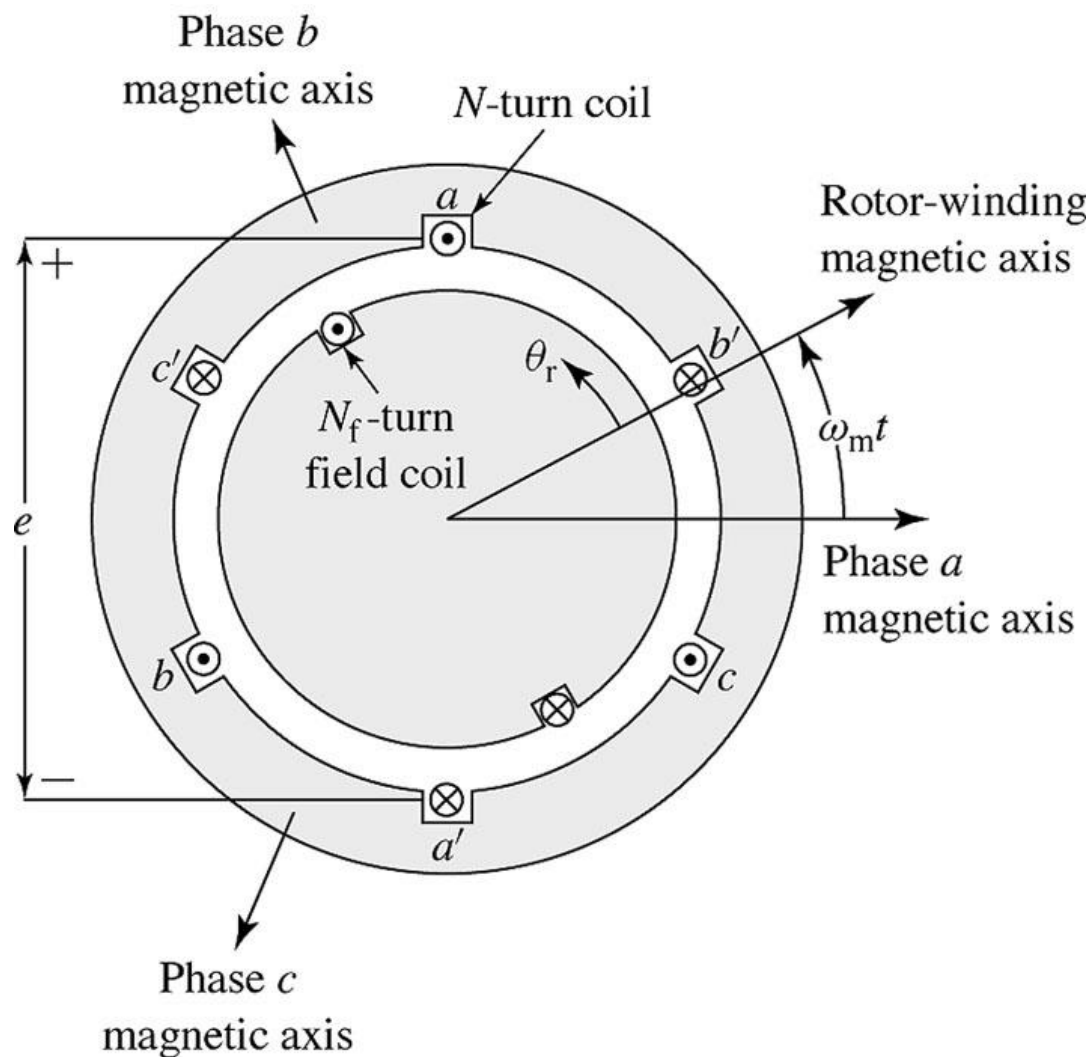


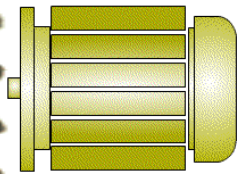


# The production of a rotating magnetic field by means of three-phase currents.

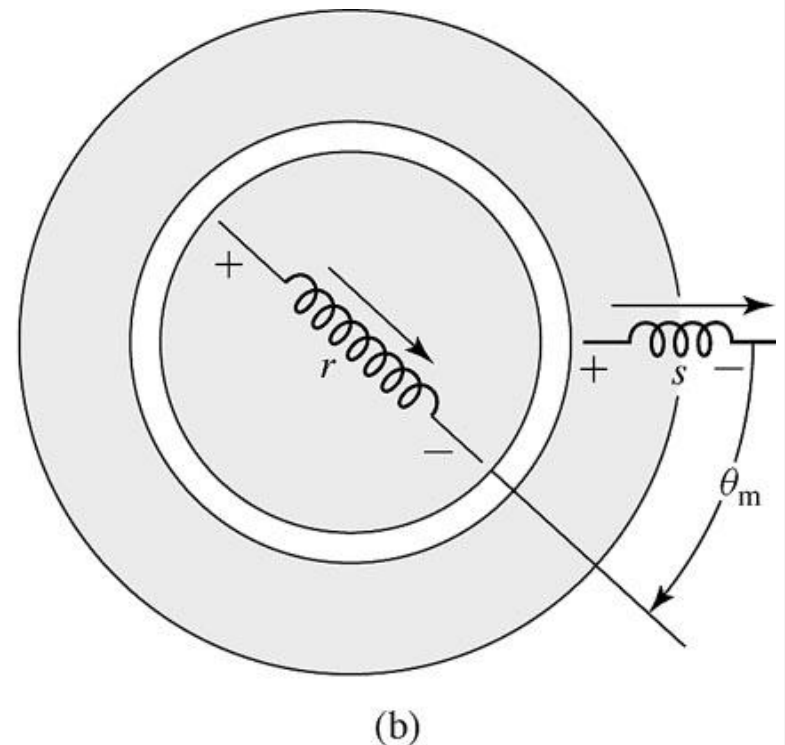
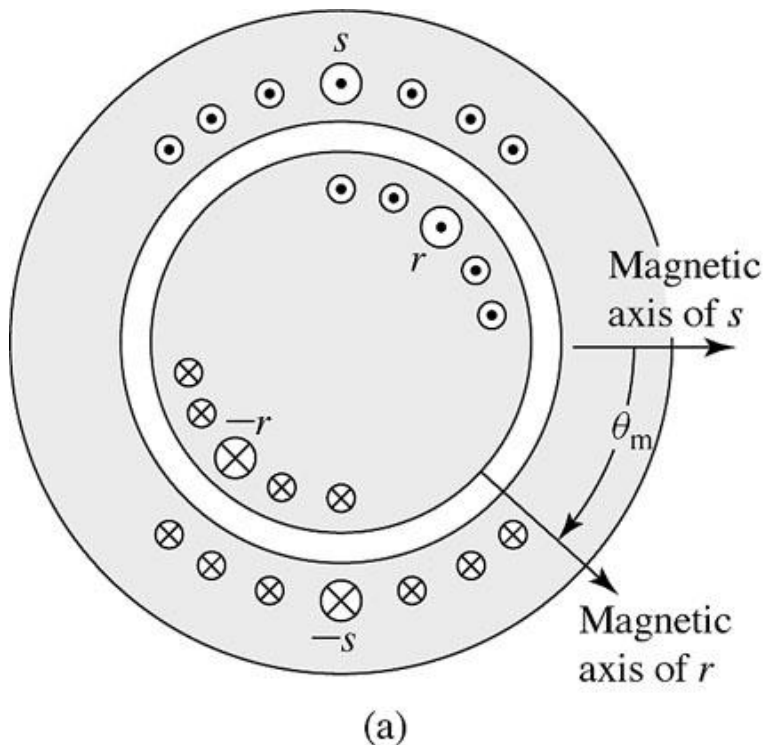
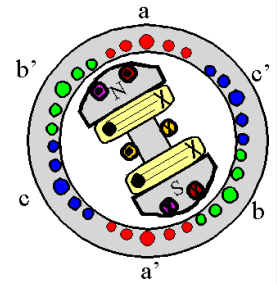


# Cross-sectional view of an elementary three-phase ac machine.

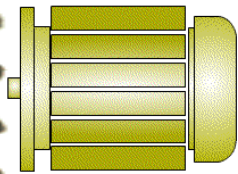




# Elementary two-pole machine with smooth air gap: (a) winding distribution and (b) schematic representation.





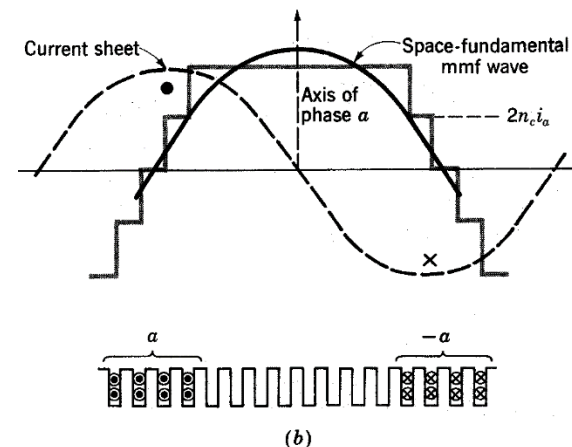
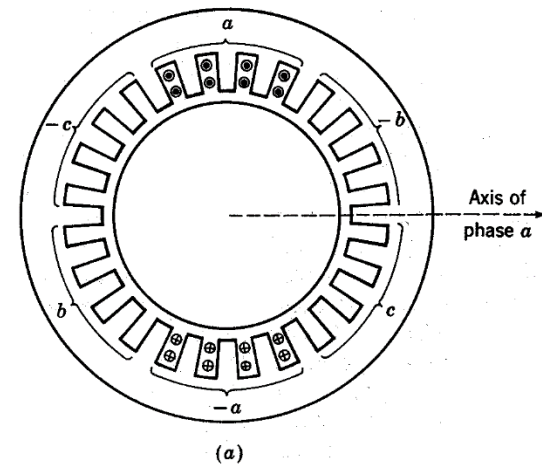
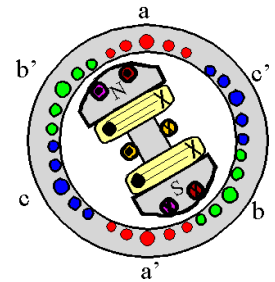


## Elementary two-pole machine with smooth air gap: cont..

- The fundamental component of the resultant *mmf* can be obtained by adding the fundamental components of these individual coils, and it can be expressed as

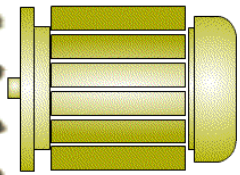
$$F_{a1} = \frac{4}{\pi} \frac{k_p N_{ph}}{P} i_a \cos \theta$$

- where  **$N_{ph}$**  is the total number of turns of the phase winding, which is formed by these coils,  **$k_d$**  is known as the distribution factor of the winding, which is defined by

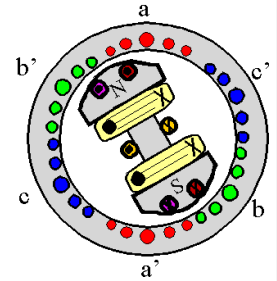


The mmf of one phase of a distributed two-pole three-phase winding with full-pitch coils.

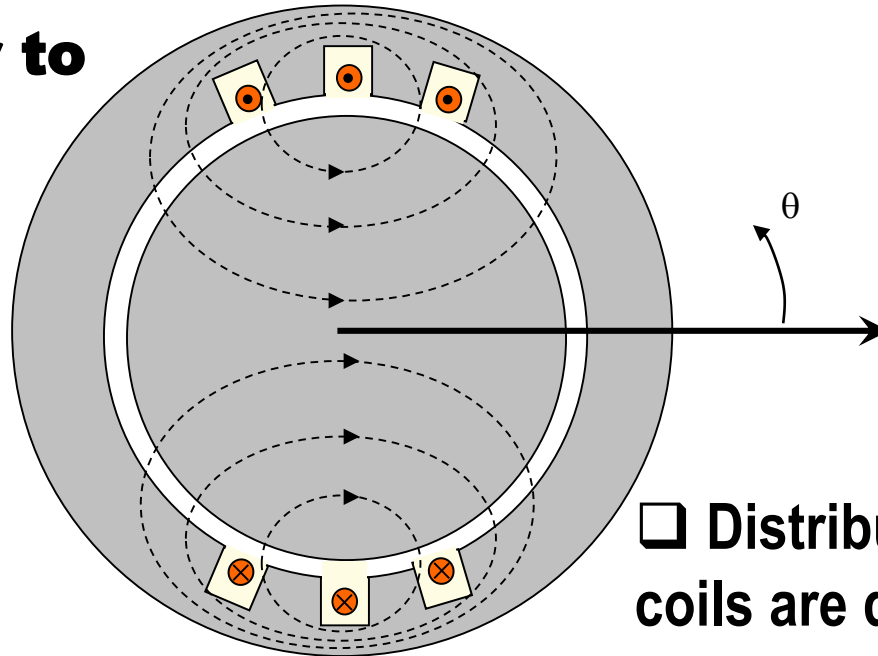




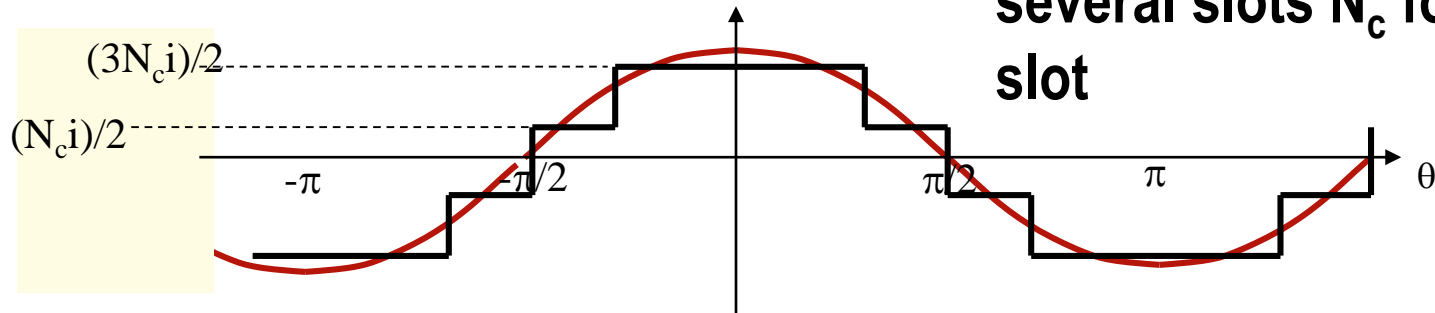
# Rotating field theory

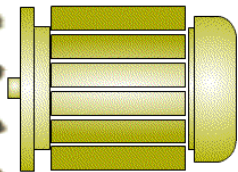


**MMF** closer to  
**sinusoidal**  
**- less**  
**harmonic**  
**contents**

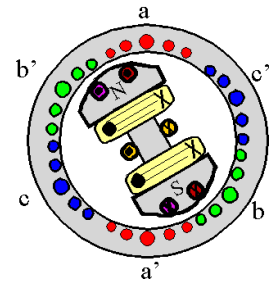


□ Distributed winding –  
coils are distributed in  
several slots  $N_c$  for each  
slot





## Rotating field theory cont..



Let  $i_a = I_m \cos \omega t$ , and we have

$$F_{a1} = \frac{4}{\pi} \frac{k_w N_{ph}}{P} I_m \cos \omega t \cos \theta$$

$$= F_m \cos \omega t \cos \theta$$

where

$$F_m = \frac{4}{\pi} \frac{k_w N_{ph}}{P} I_m$$

The *mmf* of a distributed phase winding is a function of both space and time. When plotted at different time instants as shown below, we can see that it is a pulsating sine wave. We call this type of *mmf* as a **pulsating mmf**.

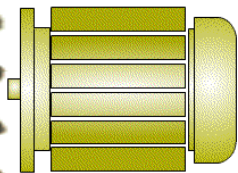
Because  $\cos \alpha \cos \beta = \frac{\cos(\alpha - \beta) + \cos(\alpha + \beta)}{2}$ , the above expression of the *mmf*

fundamental component can be further written as

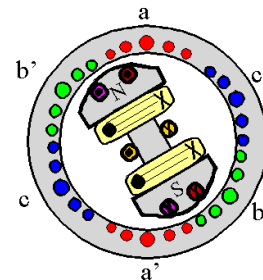
$$F_{a1} = \frac{F_m}{2} \cos(\theta - \omega t) + \frac{F_m}{2} \cos(\theta + \omega t)$$

$$= F_+ + F_-$$





# Magnetic Field of Three Phase Windings



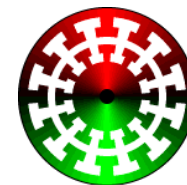
- Once we get the expression of *mmf* for a single phase winding, it is not difficult to write the expressions of *mmf*'s for three single phase windings placed **120Deg** (electrical) apart and excited by balanced three phase currents:

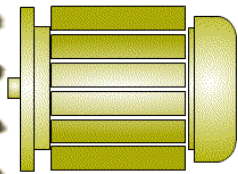
$$F_{a1} = F_m \cos \omega t \cos \theta = \frac{F_m}{2} \cos(\theta - \omega t) + \frac{F_m}{2} \cos(\theta + \omega t)$$

$$\begin{aligned} F_{b1} &= F_m \cos(\omega t - 120^\circ) \cos(\theta - 120^\circ) \\ &= \frac{F_m}{2} \cos(\theta - \omega t) + \frac{F_m}{2} \cos(\theta + \omega t - 240^\circ) \end{aligned}$$

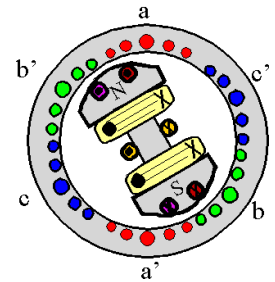
$$\begin{aligned} F_{c1} &= F_m \cos(\omega t - 240^\circ) \cos(\theta - 240^\circ) \\ &= \frac{F_m}{2} \cos(\theta - \omega t) + \frac{F_m}{2} \cos(\theta + \omega t - 480^\circ) \end{aligned}$$

and





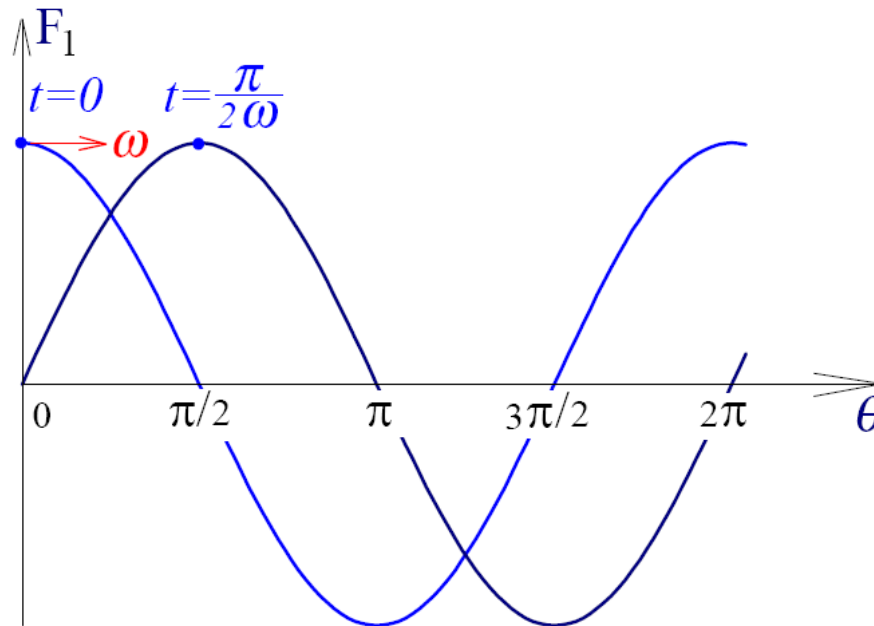
# Magnetic Field of Three Phase Windings cont...



Therefore, the resultant *mmf* generated by a three phase winding is

$$F_1 = F_{a1} + F_{b1} + F_{c1} = \frac{3F_m}{2} \cos(\theta - \omega t)$$

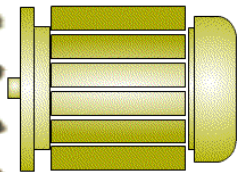
Note that  $\cos(\theta + \omega t) + \cos(\theta + \omega t - 240^\circ) + \cos(\theta + \omega t - 480^\circ) = 0$



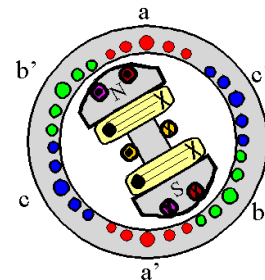
Rotating *mmf* in  $+\theta$  direction







## Magnetic Field of Three Phase Windings cont...



The above diagram plots the resultant  $mmf F_1$  at two specific time  $+\theta$  stants:  $t=0$   $a(a \rightarrow b \rightarrow c)$   $t=\pi/2\omega$ .

*It can be readily observed that  $F_1$  is a rotating mmf in the direction*

*with a constant magnitude  $3F_m/2$ . The speed of this rotating mmf can be calculated as*

$$\omega_f = \frac{d\theta}{dt} = \frac{\pi/2}{\pi/2\omega} = \omega \quad \text{rad/s (electrical)}$$

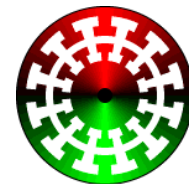
When expressed in mechanical radians per second and revolutions per minute, the speed of the rotating **mmf** can be expressed as

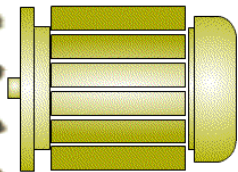
$$\omega_f = \frac{\omega}{P/2} \quad \text{rad/s (mechanical)}$$

and

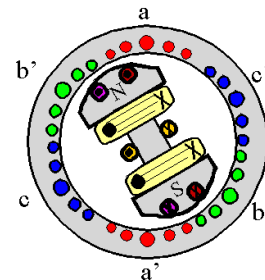
$$n_f = \frac{60\omega_f}{2\pi} = \frac{120f}{P} \quad \text{rev/min}$$

respectively

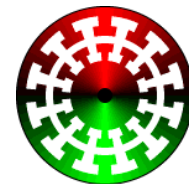


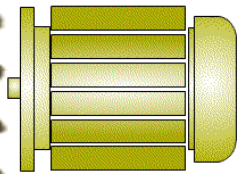


## Magnetic Field of Three Phase Windings cont...

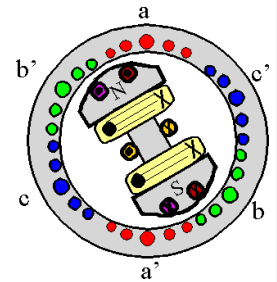


- Again, for a machine with uniform air gap, the above analysis for *mmf* is also valid for the magnetic field strength and the flux density in the air gap.
- Therefore, *the speed of a rotating magnetic field is proportional to the frequency of the three phase excitation currents, which generate the field.*





# Rotating Magnetic Field



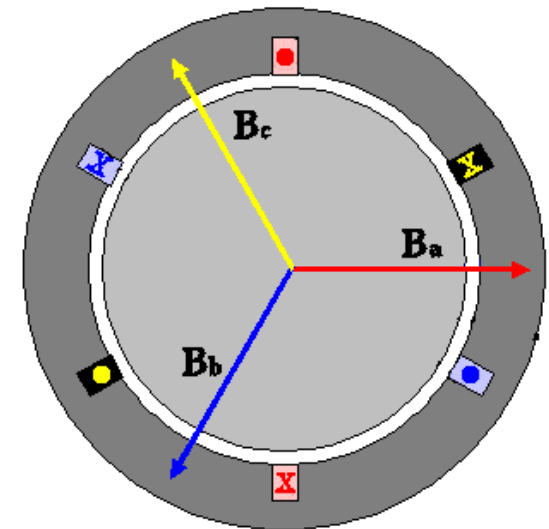
$$B_{net}(t) = B_a(t) + B_b(t) + B_c(t)$$

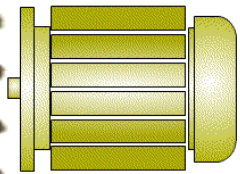
$$= B_M \sin(\omega t) \angle 0^\circ + B_M \sin(\omega t - 120^\circ) \angle 120^\circ + B_M \sin(\omega t - 240^\circ) \angle 240^\circ$$

$$= B_M \sin(\omega t) \hat{x}$$

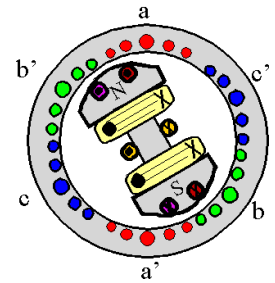
$$- [0.5 B_M \sin(\omega t - 120^\circ)] \hat{x} - \left[ \frac{\sqrt{3}}{2} B_M \sin(\omega t - 120^\circ) \right] \hat{y}$$

$$- [0.5 B_M \sin(\omega t - 240^\circ)] \hat{x} + \left[ \frac{\sqrt{3}}{2} B_M \sin(\omega t - 240^\circ) \right] \hat{y}$$

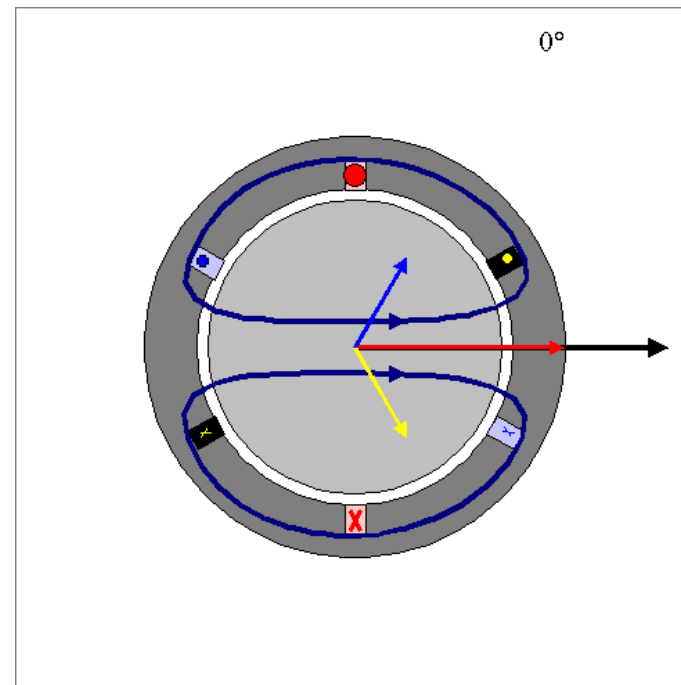


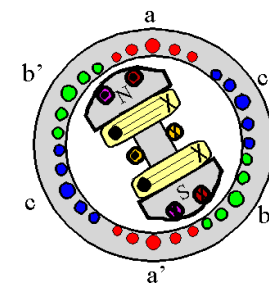
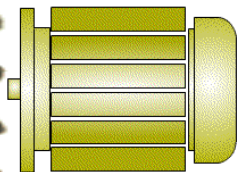


# Rotating Magnetic Field

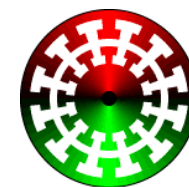


$$\begin{aligned}
 B_{net}(t) = & \left[ B_M \sin(\omega t) + \frac{1}{4} B_M \sin(\omega t) + \frac{\sqrt{3}}{4} B_M \cos(\omega t) + \frac{1}{4} B_M \sin(\omega t) - \frac{\sqrt{3}}{4} B_M \cos(\omega t) \right] \hat{x} \\
 & + \left[ -\frac{\sqrt{3}}{4} B_M \sin(\omega t) - \frac{3}{4} B_M \cos(\omega t) + \frac{\sqrt{3}}{4} B_M \sin(\omega t) - \frac{3}{4} B_M \cos(\omega t) \right] \hat{y} \\
 = & [1.5 B_M \sin(\omega t)] \hat{x} - [1.5 B_M \cos(\omega t)] \hat{y}
 \end{aligned}$$

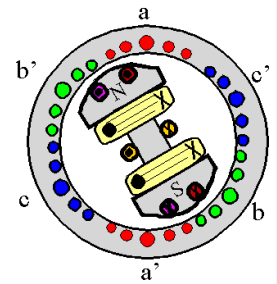
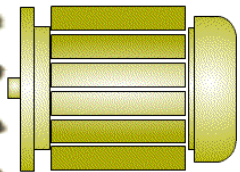




PresenterMedia







The background of the image is a spiral-bound notebook. The cover is a light beige or tan color with a fine, woven texture. A silver-colored metal spiral binding is visible along the left edge of the notebook. The text "END OF LECTURE 10" is printed in a bold, blue, sans-serif font across the center of the notebook cover.

**END OF LECTURE 10**